

S26 AerE 5740: Optimal Control

Practice Set 02: Equation-of-motion for Pendulum with Moving Pivot

Instructor: Abhishek Halder

All rights reserved

Problem 1

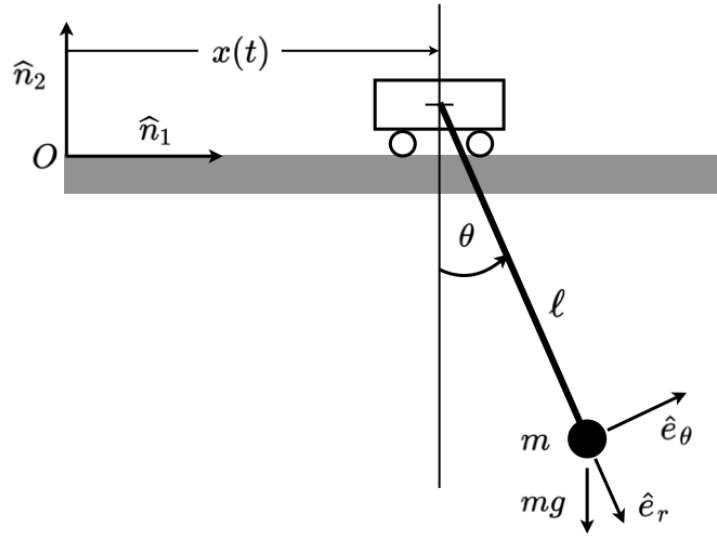


Figure 1: Pendulum with moving pivot.

Consider a pendulum with moving pivot as shown in Fig. 1 where the pivot motion is a sinusoidal function of time, i.e.,

$$x(t) = A \sin \Omega t,$$

for some given constants A, Ω . The mass of the bob is m , and the massless rod is of length ℓ , as shown. Neglect friction.

(a) Equation of motion.

Use Euler-Lagrange equation to derive the equation of motion for the pendulum.

Solution for part (a):

Since the position vector for the pendulum is $r(t) = (A \sin \Omega t) \hat{n}_1 + \ell \hat{e}_r$, we have

$$\dot{r}(t) = (A\Omega \cos \Omega t) \hat{n}_1 + \frac{d}{dt} (\ell \hat{e}_r) = (A\Omega \cos \Omega t) \hat{n}_1 + \frac{d\ell}{dt} \hat{e}_r + \dot{\theta} \hat{e}_\theta \times \ell \hat{e}_r = (A\Omega \cos \Omega t) \hat{n}_1 + 0 + \ell \dot{\theta} \hat{e}_\theta.$$

Therefore, the kinetic energy T for the pendulum is

$$\begin{aligned} T &= \frac{1}{2}m\langle\dot{r}, \dot{r}\rangle = \frac{1}{2}m \left[(A\Omega \cos \Omega t)^2 + (\ell\dot{\theta})^2 + 2\langle A\Omega \cos(\Omega t)\hat{n}_1, \ell\dot{\theta}\hat{e}_\theta \rangle \right] \\ &= \frac{1}{2}m \left[(A\Omega \cos \Omega t)^2 + (\ell\dot{\theta})^2 + 2A\ell\dot{\theta}\Omega \cos(\Omega t) \cos \theta \right]. \end{aligned}$$

The potential energy for the pendulum is $V = mg\ell(1 - \cos \theta)$. So, the Lagrangian L is

$$L = T - V = \frac{1}{2}m \left[(A\Omega \cos \Omega t)^2 + (\ell\dot{\theta})^2 + 2A\ell\dot{\theta}\Omega \cos(\Omega t) \cos \theta \right] - mg\ell(1 - \cos \theta).$$

Thus, the EL equation

$$0 = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = m\ell^2\ddot{\theta} + mg\ell \sin \theta - mA\Omega^2\ell \cos \theta \sin(\Omega t),$$

gives the equation of motion for the pendulum as

$$\ddot{\theta} + \frac{g}{\ell} \sin \theta = \left(\frac{A\Omega^2}{\ell} \right) \cos \theta \sin \Omega t.$$

(b) Conserved quantities.

(b.1) Compute the total energy E , and the Hamiltonian H . Is $H = E$ in this case?

(b.2) Use your answer in part (b.1) to conclude if any of the two quantities among H and E is a constant of motion.

Solution for part (b):

(b.1) The total energy is

$$E = T + V = \frac{1}{2}m \left[(A\Omega \cos \Omega t)^2 + (\ell\dot{\theta})^2 + 2A\ell\dot{\theta}\Omega \cos(\Omega t) \cos \theta \right] + mg\ell(1 - \cos \theta).$$

The Hamiltonian is

$$H = \frac{\partial L}{\partial \dot{\theta}}\dot{\theta} - L = -\frac{1}{2}mA^2\Omega^2 \cos^2(\Omega t) + mg\ell(1 - \cos \theta) + \frac{1}{2}m\ell^2\dot{\theta}^2.$$

Comparing the above two expressions, $H \neq E$ in this case.

(b.2) By direct differentiation, we get

$$\frac{dE}{dt} = -mA\Omega \cos(\Omega t) \left(A\Omega \sin(\Omega t) + \ell\dot{\theta} \sin \theta \right),$$

where we have used the equation of motion in part (a) to substitute for $\ddot{\theta}$. From above, we conclude that $\frac{dE}{dt} \neq 0$, i.e., E is not constant. In retrospect, this is intuitive since the constraint force needed to ensure $x(t) = A \sin \Omega t$ does work on the system.

Similarly, we have

$$\frac{dH}{dt} = mA\Omega^2 \sin(\Omega t) \left(A\Omega \cos(\Omega t) + \ell \dot{\theta} \cos \theta \right),$$

where again, we have used the equation of motion in part (a) to substitute for $\ddot{\theta}$. From above, we conclude that $\frac{dH}{dt} \neq 0$, i.e., H is not constant. This is expected since the Lagrangian L was an explicit function of time t .